

## Ligand-Stabilized Reduced-Dimensionality Perovskites

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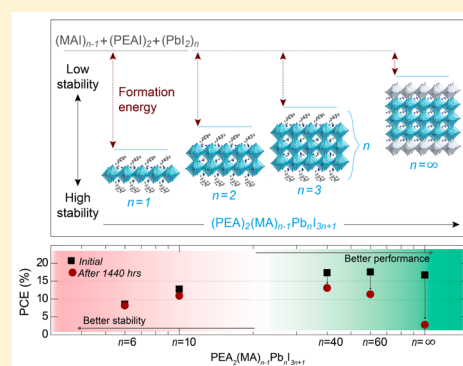
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### Supporting Information

**ABSTRACT:** Metal halide perovskites have rapidly advanced thin-film photovoltaic performance; as a result, the materials' observed instabilities urgently require a solution. Using density functional theory (DFT), we show that a low energy of formation, exacerbated in the presence of humidity, explains the propensity of perovskites to decompose back into their precursors. We find, also using DFT, that intercalation of phenylethylammonium between perovskite layers introduces quantitatively appreciable van der Waals interactions. These drive an increased formation energy and should therefore improve material stability. Here we report reduced-dimensionality (quasi-2D) perovskite films that exhibit improved stability while retaining the high performance of conventional three-dimensional perovskites. Continuous tuning of the dimensionality, as assessed using photophysical studies, is achieved by the choice of stoichiometry in materials synthesis. We achieve the first certified hysteresis-free solar power conversion in a planar perovskite solar cell, obtaining a 15.3% certified PCE, and observe greatly improved performance longevity.



## INTRODUCTION

The rapid advance of organic–inorganic metal halide perovskites represents a breakthrough for next-generation thin-film optoelectronics.<sup>1–6</sup> In particular, methylammonium (MA, CH<sub>3</sub>NH<sub>3</sub>) lead iodide (MAPbI<sub>3</sub>) perovskites possess an appropriate bandgap energy,<sup>7,8</sup> large absorption coefficient, and long-range ambipolar photocarrier diffusion.<sup>9</sup> Tremendous efforts from multiple research groups have propelled the certified power conversion efficiency (PCE) of perovskite solar cells to 20.1%.<sup>10</sup>

Unfortunately, perovskite materials lack long-term stability under ambient operation conditions. This has been ascribed, with the aid of theoretical studies, to the low formation energy of MAPbI<sub>3</sub>.<sup>11</sup> This challenge has motivated efforts to improve device stability, and progress in this direction has been made by developing novel device architectures.<sup>12</sup>

Here we focus instead on the stability of the active material itself. We investigate, using theory and experiment, the origins of instability in conventional MAPbI<sub>3</sub>, and we propose and fabricate enhanced perovskites—taking advantage of the versatile physical chemistry of this materials system—that constitute low-dimensional hybrid organic–inorganic crystalline thin films. There exist prior reports of two-dimensional (2D) layered perovskites, crafted using bulkier organic molecules in place of MA, that exhibit enhanced stability;

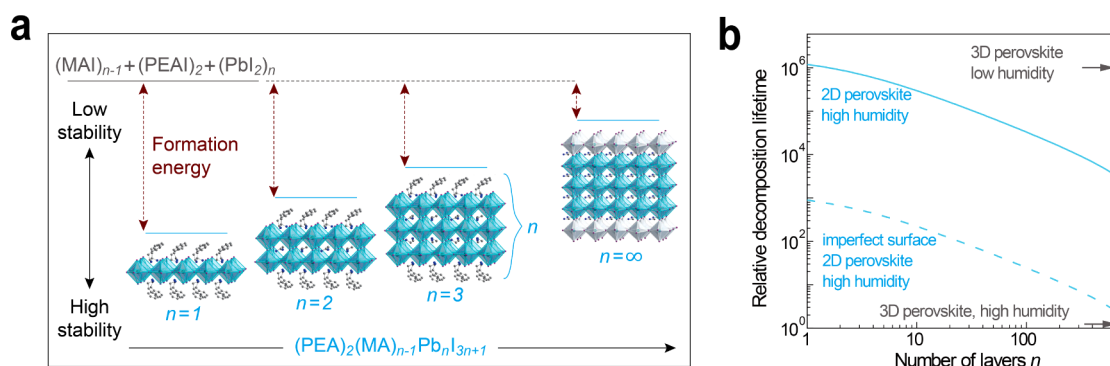
however, these suffer from poor photovoltaic performance.<sup>13</sup> These prior reports have shown PCEs lower than 5% due to their low absorption coefficient and poor carrier transport,<sup>14</sup> and this has hindered their further application. Here, we develop and investigate a platform of mixed-organic, dimensionally tuned quasi-2D perovskite thin films that continuously bridge the gap between 2D and 3D materials.<sup>15,16</sup> We find that members of the quasi-2D family combine the enhanced stability of 2D perovskites with the excellent optoelectronic parameters, including long-range photocarrier diffusion, of 3D perovskites. Our approach relies on computational methods to unravel the origins of instability in conventional MAPbI<sub>3</sub>, while complementary studies of the physical and optoelectronic properties of quasi-2D perovskites allow us to identify optimized conditions to design and fabricate chemically stable, high-efficiency photovoltaic devices.

## RESULTS AND DISCUSSION

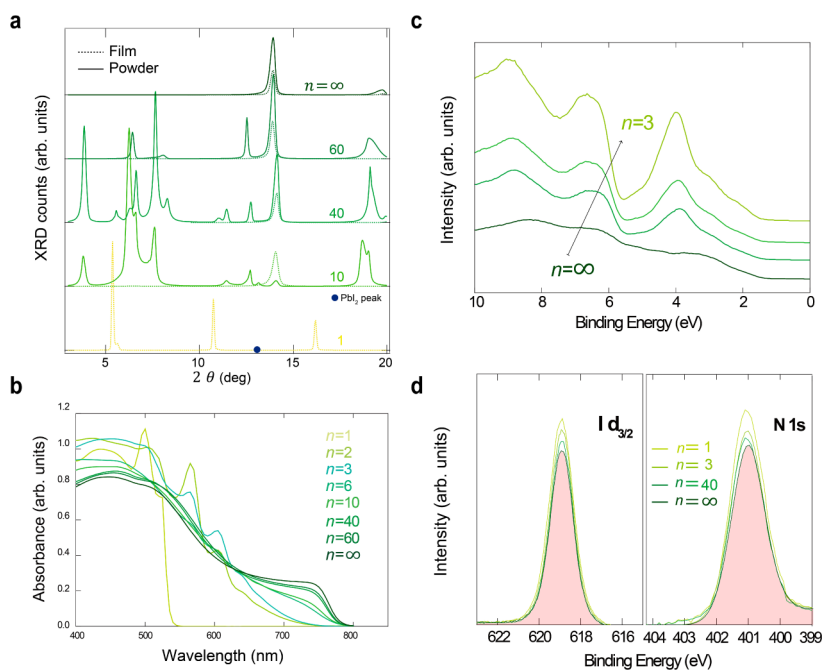
The intermediates between 3D and 2D are layered perovskites systematically synthesized by introducing a large organic cation, phenylethylammonium (PEA, C<sub>8</sub>H<sub>9</sub>NH<sub>3</sub>), at a judiciously chosen stoichiometry. This cation possesses a large ionic radius

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**Figure 1.** Energetics of perovskite formation and stability. (a) Unit cell structure of  $(\text{C}_8\text{H}_9\text{NH}_3)_2(\text{CH}_3\text{NH}_3)_{n-1}\text{Pb}_n\text{I}_{3n+1}$  perovskites with different  $n$  values, showing the evolution of dimensionality from 2D ( $n = 1$ ) to 3D ( $n = \infty$ ). (b) DFT simulation of the formation energy of perovskite with different  $n$  values in different atmospheres.



**Figure 2.** Structure and optical characterization of dimensionally tuned perovskite films. (a) Low-diffraction-angle region of XRD spectra. (b) Absorption spectra of the perovskite films with different  $n$  values. (c) Valence band photoemission. (d) XPS data for the atomic core-level peaks.

which is incompatible with a 3D perovskite structure; therefore, perovskites based on this cation tend to crystallize into a layered, 2D structure.<sup>17</sup> Consequently, we can continuously tune the dimensionality of metal halide perovskite compounds by mixing stoichiometric quantities of lead iodide ( $\text{PbI}_2$ ), methylammonium iodide (MAI,  $\text{CH}_3\text{NH}_3\text{I}$ ), and phenylethylammonium iodide (PEAI) to yield compounds with different layer ( $n$ ) values in the series  $\text{PEA}_2(\text{CH}_3\text{NH}_3)_{n-1}\text{Pb}_n\text{I}_{3n+1}$  (Table S1). In this notation, the limit  $n = \infty$  corresponds to the cubic 3D perovskite,  $\text{CH}_3\text{NH}_3\text{PbI}_3$ , while the other  $n$  values describe 2D ( $n = 1$ ) or quasi-2D ( $n > 1$ ) perovskite structures (Figure 1a).

We investigated the stability of dimensionally tuned perovskites using density functional theory (DFT) simulations. Calculations show that the conventional 3D perovskite is marginally stable relative to decomposition into bulk  $\text{PbI}_2$  and MAI constituents (see details for DFT simulations in Supporting Information, Figure S1), and indeed both components have been observed in thermal gravimetric analysis.<sup>18</sup> However, this process cannot happen within the

bulk spontaneously and needs to start from the surface, where molecules have sufficient freedom to rearrange. We further analyze the formation enthalpy into that of constituent layers, predicting  $\sim 0.55$  eV per molecule required to recrystallize the MAI layer and  $\sim 0.6$  eV gain upon conversion of the  $\text{PbI}_2$  layer (see details on DFT simulations and results and discussion in the Supporting Information). Thus, MAI serves as a protection layer, limiting perovskite decomposition, while the  $\text{PbI}_2$  surface termination is inherently unstable.

Under ambient conditions, the degradation of 3D perovskites is known to result in the formation of a  $\text{PbI}_2$  yellow phase, and X-ray photoelectron spectroscopy (XPS) suggests that it is accompanied by the loss of MAI into the gas phase, with no measurable formation of perovskite hydrate,<sup>19</sup> even in a high-humidity environment. In contrast, simulations predict a greater than 2 eV desorption energy for organic molecules. To reconcile these observations, we included water in our simulations<sup>20</sup> and found a significant decrease in desorption energies and decomposition lifetimes (Figure 1b), consistent with the sensitivity of perovskite to moisture.<sup>18,33</sup>

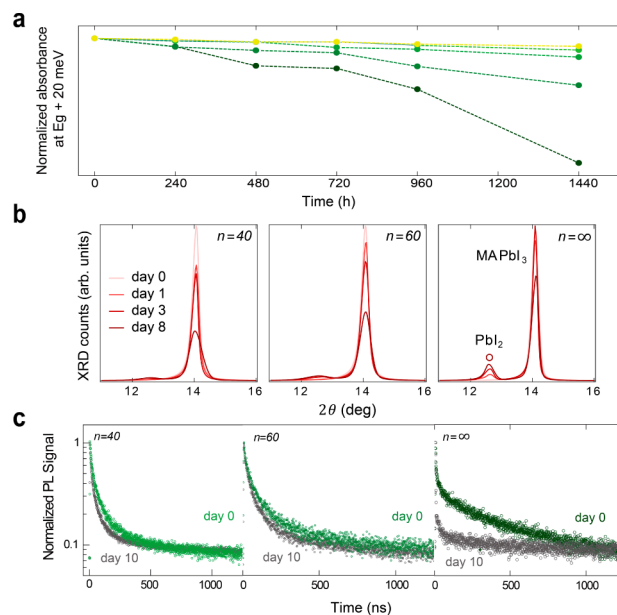
Notably, we find that van der Waals interactions between the capping organic molecules are key drivers of increased material stability.<sup>21</sup> The energy required to remove PEA from the perovskite is 0.36 eV higher than that for MAI, reducing the desorption rate by 6 orders of magnitude, and slowing film decomposition 1000-fold (additional details on DFT simulations and results can be found in Supporting Information).

Motivated by these insights, and building on the previously reported solvent-engineering fabrication process,<sup>22</sup> we fabricated dimensionally tuned perovskite films using single-step spin-coating. X-ray diffraction (XRD) patterns for the 3D and quasi-2D perovskite thin films show very similar reflections, indicating a preferential orientation of the crystallites relative to the substrate. Thus, thin-film samples do not exhibit Bragg reflections at low diffraction angles  $2\theta$ , yet these low-diffraction-angle features represent the unambiguous signatures of the layered structure, specifically of the elongated unit cell.<sup>23</sup>

We therefore obtained powders from the films that would lack preferential orientation. These powder samples show diffraction scans that, in all but the 3D case ( $n = \infty$ ), exhibit a series of Bragg peaks at low angles ( $2\theta < 10^\circ$ ).<sup>24</sup> The XRD data confirm the formation of lead iodide  $n$ -layer perovskite sheets (Figure 2a) but are not conclusive on the phase purity of quasi-2D films. Consequently, we embarked on a study of the phase composition of dimensionally tuned perovskite solids of various  $n$ .

As expected, decreasing the dimensionality of the perovskite structure causes an increase in the electronic bandgap.<sup>30,31</sup> The dimensional tuning of the electronic bandgap,  $E_G$ , can be followed from the absorption spectra (Figure 2b), which show a higher-energy absorption edge for lower  $n$  values. The trend is in good agreement with DFT calculations (Figure S3). The energy levels of these compounds have been analyzed using a combination of ultraviolet photoemission spectroscopy (UPS) and UV-vis absorption spectra, as shown in Figure 2c. The valence band maximum (VBM) for 3D  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite is found 5.34 eV below the vacuum level. From this and the optical bandgap, we conclude that the conduction band minimum (CBM) is located at 3.77 eV.<sup>25</sup> When we decrease  $n$  values, the CBM becomes shallower, while the VBM shows a negligible variation (Figure 2c). The stoichiometry of these perovskite films was verified using XPS (Figure 2d), which revealed a gradual decrease of both the I:Pb and N:Pb ratios throughout the film with increasing  $n$  values, in good agreement with the values expected from the nominal molar ratio of precursors (Table S1). The above studies, taken together, indicate a controlled synthesis of mixed-stoichiometry materials in which the organic cation enables continuous tuning of the dimensionality, and resultant optical properties, of the perovskite compounds.

We then investigated material stability as a function of dimensionality. The perovskite films were stored for extended periods of time in environments having controlled humidity levels. In the case of bulk  $\text{MAPbI}_3$  perovskites, the absorbance (Figure 3a) near the absorption edge onset degrades rapidly over two months when the sample is stored in 90% relative humidity (RH). After two months, the absorbance of 3D perovskite films at wavelengths longer than 500 nm is significantly suppressed, signaling full reversion to  $\text{PbI}_2$ . In contrast, the pure 2D perovskite ( $n = 1$ ),  $\text{PEA}_2\text{PbI}_4$ , exhibits excellent stability over the same period of time using the same measures, with no ostensible variation in the absorbance observed, which is in good agreement with previous reports.<sup>16</sup>



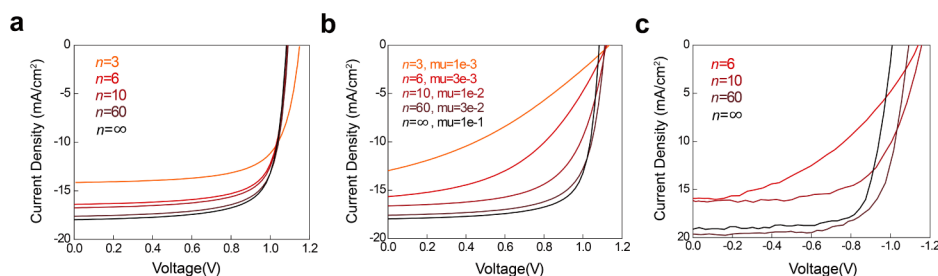
**Figure 3.** Photophysical and structural stability of 3D and quasi-2D perovskite films. (a) Relative absorption intensity near the absorption onset for different  $n$  values. (b) Film XRD spectra as a function of  $n$  at high humidity level (90%) with different aging time. (c) PL decay data before and after 10 days storage.

As expected, the quasi-2D perovskites ( $n = 10, 40, 60$ ) exhibit significantly improved stability compared to 3D perovskite. A slight change in the absorbance is observed for the prominent perovskite peak around 750 nm, suggesting that a part of these perovskite films has converted to  $\text{PbI}_2$  (Figure 3a). Layered perovskites with lower  $n$  value ( $n = 10$ ) (i.e., closer to 2D) are found to exhibit better stability of optical properties compared to those with higher  $n$  values ( $n = 40, 60$ ).

We tracked the materials' structure stability using XRD periodically recorded on perovskite films stored at 90% RH for more than a week (Figure 3b). Strong (002) diffraction peaks (at  $2\theta = 13.98^\circ$ ) characteristic of the perovskite structure are observed for all fresh samples. However, a new peak at  $2\theta = 12.42^\circ$  emerges for 3D perovskite film, corresponding to a Bragg reflection of  $\text{PbI}_2$  structure, with an increasing intensity over time. In contrast, only a very weak  $\text{PbI}_2$  reflection appears in the  $n = 60$  and  $n = 40$  perovskite films, with a peak intensity nearly unchanged over the same period of time, which indicates a significant improvement of stability compared to 3D perovskite.

We also investigated the materials' stability using transient photoluminescence (PL) decay measurements. We then tracked any changes over time (Figure 3c). Dramatic changes are observed in the PL decay traces for the 3D perovskite after 10 days storage in air. There is no observable variation in the carrier lifetime of quasi-2D ( $n = 40, 60$ ) perovskites. The PL decay changes provide further support of enhanced stability: specifically, the lower  $n$ -valued perovskites exhibit improved stability of PL lifetime compared to higher  $n$ -values counterparts, a finding that is in good agreement with DFT studies (Figure 1b).

We turned to evaluating the enhanced-stability materials for photovoltaic applications. Unfortunately, the pure 2D perovskites ( $n = 1$ ) and low- $n$ -value quasi-2D perovskite films exhibit appreciably reduced mobility ( $\mu$ ) and lower lifetime;<sup>21</sup> further, they have a significantly wider bandgap that will sacrifice

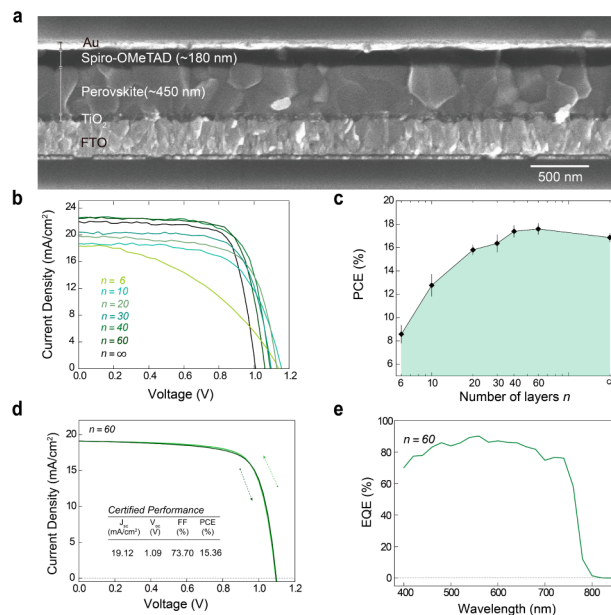


**Figure 4.** Comparison of  $J$ - $V$  curves for different numbers of perovskite layers. (a) Simulation using experimental absorption spectra and band gaps, and mobility of  $0.1 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ . (b) Simulation with radiative recombination of  $5 \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$  and gradually decreasing mobility for lower  $n$ . Same trap density is used for all cases. Perovskite layer was estimated to be 200 nm thick. (c) Experimental curves.

absorption in the near-infrared (NIR) portion of the solar spectrum. In contrast, our studies of carrier lifetime ( $\tau$ ) and carrier diffusion length ( $L_D$ )<sup>9</sup> (Table S2) indicated that the quasi-2D perovskite films having higher  $n$  values ( $n = 40, 60$ ) exhibit impressive  $\tau$ ,  $\mu$ , and  $L_D$  compatible with photovoltaic performance.<sup>27</sup> Since the 3D perovskite films are limited not by mobility but instead by trap volume density,<sup>26</sup> we reasoned that some sacrifice to mobility in the quasi-2D materials may potentially be tolerable before radiative recombination becomes a considerable loss, an idea confirmed via self-consistent optoelectronic simulations. Figure 4 shows a qualitative comparison of simulated and experimental  $J$ - $V$  curves. Modeling assuming the same carrier diffusion length of 350 nm for all samples and using experimental absorption spectra and band gaps (Figure 4a) cannot reproduce experimental trends (Figure 4c): (i) the fill factor (FF) stays the same for all samples, and for low  $n$  it is higher than experimental values; (ii)  $V_{OC}$  is practically unaltered, following the band gap trends, while experimental  $V_{OC}$  is higher for lower  $n$ ; (iii) the changes in absorption are not sufficient to explain the drop of  $J_{SC}$  for lower  $n$ . The FF is primarily defined by the carrier diffusion length, which in turn depends on trap density. However, introducing more traps will lead to deterioration of all parameters,  $J_{SC}$ , FF, and  $V_{OC}$ , while experimentally  $V_{OC}$  actually increases for lower  $n$ . The increase of  $V_{OC}$  is usually a signature of decreased mobility, since slower carrier extraction allows for more carrier accumulation, and thus, higher  $V_{OC}$ . Higher carrier accumulation, in turn, increases the chances for carriers to meet and recombine radiatively. Including the radiative recombination, in combination with lower mobility, allows reproducing experimental trends, i.e., decrease in  $J_{SC}$  and FF while increasing  $V_{OC}$  (Figure 4b). In sum, weak absorption and higher bandgap are detrimental only for very-low- $n$  devices; for  $n = 6$  and higher, the main limiting factor is the reduced mobility, leading to strong radiative recombination losses. This suggests that quasi-2D perovskites, with high  $n$  values, may be capable of performance comparable to that of 3D materials with increased stability.

We therefore constructed solar cells employing the various active layers in a planar heterojunction architecture (Figure 5a). Cross-sectional SEM indicates the presence of two uniform layers, a  $\text{TiO}_2$  layer ( $\sim 50 \text{ nm}$ ) coated with a dense perovskite absorber ( $\sim 450 \text{ nm}$ ). For the absorber layer, a dense-grained uniform morphology with grain sizes in the range of 100–500 nm was observed in all samples. The grain size shows a slight reduction with decreasing  $n$  value (Figure S4).

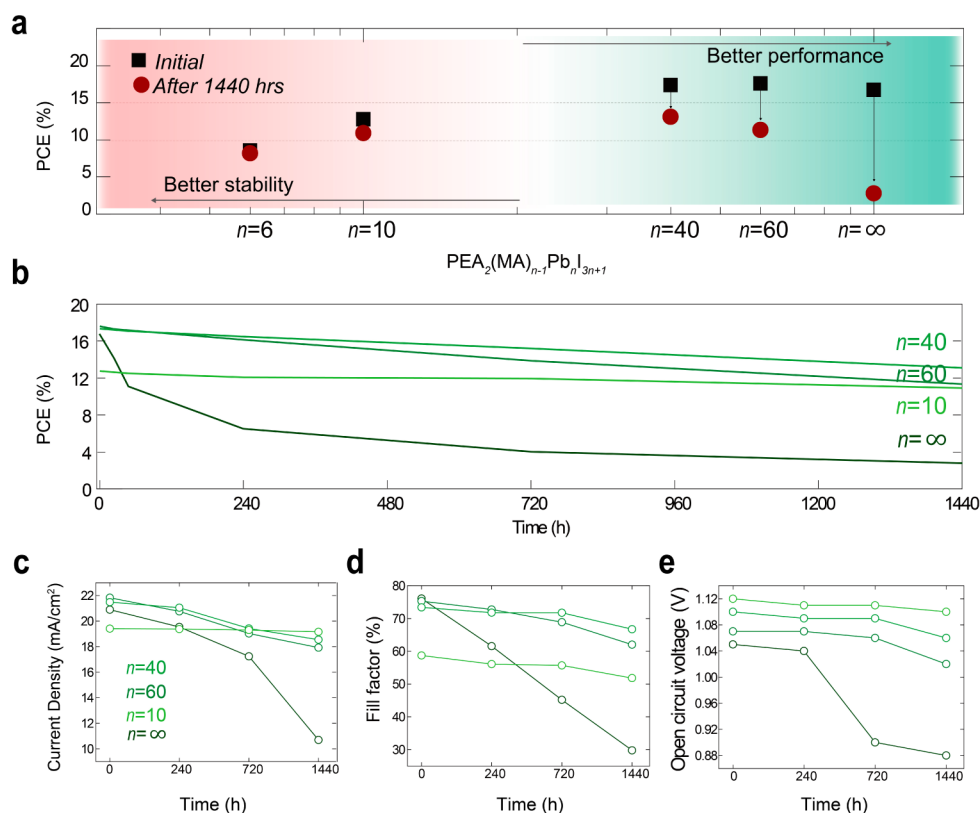
The perovskite with  $n < 40$  shows lower performance due to lower mobility, which leads to higher charge accumulation and increased radiative recombination losses, which negatively



**Figure 5.** Solar cell device architecture and performance. (a) Cross-sectional scanning electron microscopy (SEM) image for a quasi-2D perovskite device with  $n = 60$ . (b,c)  $J$ - $V$  curve and PCEs for champion perovskite device with different  $n$  values. (d,e)  $J$ - $V$  curve, both forward and reverse scan, and EQE spectra of the certified device with  $n = 60$  perovskite.

impact the FF (Figure 5b,c). In addition, perovskite films with dimensionality approaching the 2D limit ( $n < 10$ ) have a wider bandgap and reduced absorption across the solar spectrum. As expected, devices based on perovskite active layers showing a very low  $n$  value ( $n = 6$ ) exhibited very low initial performance in view of their low absorbed photon current and poor carrier transport (Figures 2b and 5b).

The top-performing quasi-2D perovskite ( $n = 60$ ) devices exhibited a hysteresis-free certified AM1.5 PCE of 15.3% (Figure 5d,e). Certification was obtained by an independent photovoltaic calibration laboratory (Figure S5). The steady-state current density measurement was performed for the 3D and  $n = 60$  perovskite devices (Figure S6). Photocurrent decay during testing of devices was observed in the case of 3D perovskites, while devices with  $n = 60$  perovskite showed constant current density over time, consistent with quasi-2D perovskites' low hysteresis. To the best of our knowledge, this is the first report of certified, hysteresis-free planar perovskite solar cells. The introduction of these multilayered perovskite structures might alter the mechanisms of electronic and ionic motion across the film, thus promoting the suppression of the



**Figure 6.** Solar cell device stability. (a) Device performance as a function of  $n$  value, which shows that increased performance was achieved with increased  $n$  value; however, in the meantime, stability was decreased. (b) Evolution of device performance as a function of time with different  $n$  values. Record of current density (c), fill factor (d), and open circuit voltage (e) change for different perovskite devices stored in nitrogen over two months.

well-known hysteresis in the current–voltage characteristic of planar 3D perovskite solar cells.<sup>28</sup>

We measured solar cell performance and its evolution in time (Figure 6a). The 3D perovskite devices began with 16.6% PCE but degraded over 8 weeks to less than 3% PCE when stored in  $N_2$ . Moderate  $n$  value perovskites ( $n = 10$ ) showed improved performance together with good long-term stability (Figure 6b–e); however, the absolute device performance is still low compared with that of the state-of-the-art perovskite device due to poor carrier transport, which limits their further application.

Most interesting were the quasi-2D perovskites, both  $n = 60$  and  $n = 40$  (Figure 6b–e). The performance of the  $n = 60$  device declined to 11.3% PCE after 60 days under a low-humidity atmosphere, while the  $n = 40$  device had a similar initial PCE and declined to 13.1% after 60 days (Tables S4–S8). Longer-term stabilization is expected when a physical moisture barrier such as a carbon electrode is deployed into the device architecture, combined with use of quasi-2D perovskites ( $n = 60, 40$ ).<sup>12</sup>

In analogous studies under humid air (55% RH), the 3D perovskite lost the majority of its performance, decreasing to 4.2% PCE in 3 days. In contrast, over 2 weeks under high RH, the  $n = 60$  and  $n = 40$  devices decreased to about 13% PCE in each case (Table S8). Statistical data on the decay of performance are summarized in the Supporting Information (Figure S7). The thermal stability of the devices was evaluated by annealing at 100 °C for 36 h (Figure S8). The PCE of the device with  $n = 60$  perovskite showed a modest improvement in thermal stability compared with the 3D perovskite devices.

## CONCLUSION

In sum, the very low formation energy of 3D perovskites accounts for their low stability relative to their precursor phases. This problem is exacerbated by water and humidity. We show herein that addition of an organic molecule that does not incorporate into the 3D lattice can be used to produce capped-layered quasi-2D perovskite-based materials, and that these possess much higher formation energies, owing to appreciable van der Waals forces that confer improved stability. By replacing the methylammonium with the formamidinium ( $\text{HN}=\text{CHNH}_3^+$ ) cation to further reduce quasi-2D perovskites' bandgap, solar devices with even higher performance without hysteresis could potentially be engineered in the future.<sup>10</sup> The present work suggests that, with further progress, the benefits of perovskites—in long diffusion length for high current and fill factor, as well as low trap state density for high voltage—may be incorporated into more robust solar cells.

## EXPERIMENTAL SECTION

**Phenylethylammonium Iodide Synthesis.** Phenylethylamine (10 mL, 80 mmol) was dissolved in 10 mL of ethanol in an ice bath, and HI aqueous solution (55%, 20 mL, 170 mmol) was added dropwise to the flask with vigorous stirring. After the addition of HI, the mixture was allowed to stand in ice bath for another 20 min until a colorless precipitate appeared. The precipitate was filtered under vacuum and washed thoroughly using cold diethyl ether. The crude precipitate was collected and recrystallized in ethanol twice to obtain thin, plate-like white crystals (16.5 g, yield 83%). EI-MS:  $m/z$  249.2, calcd 249.0.

**Film Fabrication.** The series of different-dimensional perovskite  $\text{PEA}_2(\text{CH}_3\text{NH}_3)_{n-1}\text{Pb}_n\text{I}_{3n+1}$  solutions were prepared by dissolving

specific stoichiometric quantities of  $\text{PbI}_2$ , MAI, and PEAI in DMF ( $n < 10$ ) or  $\gamma$ -butyrolactone/DMSO (5:5, volume ratio) mixture ( $n \geq 10$ ) at 60 °C for 12 h with vigorous stirring. The resulting solution was filtered by PTFE syringe filter (0.2  $\mu\text{m}$ ) before deposition. The DMF solution ( $n < 10$ ) was then deposited onto the substrate by a single-step spin-coating process at 5000 rpm for 30 s. The mixed  $\gamma$ -butyrolactone/DMSO solution ( $n \geq 10$ ) was coated onto the substrate via a consecutive two-step spin-coating process at 1000 and 5000 rpm for 10 and 60 s, respectively. During the second spin step, 100  $\mu\text{L}$  of chlorobenzene was deposited onto the substrate. The resulting films were then annealed at 100 °C for 1 h for better crystallization.

**Photovoltaic Device Fabrication.** A dense compact  $\text{TiO}_2$  electron transport layer (50 nm) was deposited onto a F-doped  $\text{SnO}_2$  (FTO, Pilkington, TEC 15) substrate, followed by drying at 130 °C for 10 min and calcining at 500 °C for 1 h. The  $\text{TiO}_2$  precursor solution was prepared by dissolving 890 mg of titanium isopropoxide (TTIP) and 30 mg of HCl (37%) solution in 8 mL of ethanol. The mixture was then stirred at room temperature overnight before use. The resulting  $\text{TiO}_2$  film was immersed in 120 mM  $\text{TiCl}_4$  aqueous solution at 70 °C for 1 h and then calcined at 500 °C for another 30 min. The perovskite absorber was then deposited onto the  $\text{TiO}_2$  using the method described before. The hole-transport layer was spin-coated onto the perovskite film at 4000 rpm for 30 s using a chlorobenzene solution that contained of 63 mg/mL of spiro-OMeTAD and 20  $\mu\text{L}$  of *tert*-butylpyridine, as well as 70  $\mu\text{L}$  of bis(trifluoromethane)sulfonimide lithium salt (170 mg/mL in acetonitrile). Top electrodes were deposited using an Angstrom Engineering deposition system in an Innovative Technology glovebox. The contacts consisted of 100 nm electron-beam-deposited gold, deposited at 0.4 Å/s.

**Relative Humidity Control.** For all characterizations, the perovskite films were stored at room temperature (measured as  $25 \pm 1$  °C) in a controlled-humidity desiccator. The relative humidity (RH = 90%) was achieved with an accuracy of  $\pm 3\%$  by a water/glycerol solution with careful adjustment of the weight ratio of glycerol. The precise RH was measured periodically using a calibrated hygrometer.

**Photovoltaic Performance Characterization.** The steady-state open-circuit voltage,  $V_{\text{OC}}$ , was first measured using a Keithley 2400 instrument by fixing the current to zero and sampling the voltage at multiple time points. The steady-state short-circuit current,  $J_{\text{SC}}$ , was measured by setting the bias voltage to zero and sampling the current at multiple time points. Instantaneous  $J$ - $V$  curves were then measured with a scanning rate of 0.02 V  $\text{s}^{-1}$ . The maximum power point ( $V_{\text{MPP}}$ ) was determined from the instantaneous  $J$ - $V$  curve. The steady-state PCE was measured by setting the bias voltage to the estimated  $V_{\text{MPP}}$ . Under the bias of  $V_{\text{MPP}}$ , the current density value was sampled during 15 s to obtain  $J_{\text{MPP}}$ . The PCE was finally obtained as the product of  $V_{\text{MPP}}$  and  $J_{\text{MPP}}$ . The active area was determined by the aperture before the solar cell to avoid overestimating the photocurrent. Through this aperture (area 0.049  $\text{cm}^2$ ), illumination intensity was calibrated using a Melles-Griot broadband power meter and set to be 1 sun (100 mW  $\text{cm}^{-2}$ ). The AM 1.5 solar power is supplied by a class A (<25% spectral mismatch) solar simulator (Science Tech). The spectral mismatch of the system was characterized using a calibrated reference solar cell (Newport). The accuracy of the power measurement was estimated to be  $\pm 5\%$ .

**UPS Measurements.** UPS measurements were carried out in an ultra-high-vacuum (UHV) chamber with a base pressure below  $5 \times 10^{-9}$  mbar. The photon line width was  $\sim 250$  meV and the minimum spot size  $\sim 1$  mm. He I photons (21.2 eV) were used to acquire the spectra at normal emission. The photoelectrons were collected by using a SPHERA U7 hemispherical energy analyzer with a 7-channel MCD detector, in constant analyzer energy mode. The binding energy values shown with 10 meV precision should be rounded to the nearest 100 meV value in accordance with the overall energy resolution.

**XPS Measurements.** XPS measurements were carried out in a Thermo Scientific K-Alpha system, with a 300  $\mu\text{m}$  spot size, 75 eV pass energy, and energy steps of 0.05 eV. All signals are normalized to Pb for direct comparison between different samples.

**XRD Measurements.** XRD measurements on oriented films were conducted on a Panalytical X'Pert Pro diffractometer with a Bragg-Brentano geometry and PIXCEL 1D detector equipped with a nickel filter. Unoriented scans were collected on powders loaded into a 0.1 mm diameter silicon capillary on a Bruker  $\text{D}_8$  Venture diffractometer using monochromated  $\text{Mo K}\alpha_1$  (0.71073 Å) radiation.

**Other Characterizations.** The external quantum efficiency (EQE) spectrum for photovoltaic device was obtained using a previously published process.<sup>32</sup> Samples were tested in  $\text{N}_2$  ambient. Photoluminescence (PL) was measured using a Horiba Fluorolog time-correlated single-photon-counting system with photomultiplier tube detectors. For steady-state measurements, the excitation source was a monochromatized Xe lamp. For time-resolved PL, the excitation source was a laser diode with a wavelength  $\lambda = 375$  nm; the combined (source + instrumental response function) time resolution was  $\Delta t \approx 0.13$  ns.

## ■ ASSOCIATED CONTENT

### 📄 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.5b11740.

Materials and methods, additional results and discussion, including Figures S1–8 and Tables S1–S8, and references (PDF)

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### Notes

The authors declare no competing financial interest.

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